

GEODESIC MAPPINGS OF WEAKLY BERWALD SPACES  
AND BERWALD SPACES ONTO RIEMANNIAN SPACES

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**Abstract:** The present paper is devoted to the investigation of the geodesic mappings of weakly Berwald spaces and Berwald spaces which are special Finsler spaces  $F_n$ , onto the Riemannian spaces  $\bar{V}_n$ . We can see that not every Finsler space admits non-trivial geodesic mapping onto a Riemannian space.

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1. Geodesic Mappings of Finsler Spaces

It is known [4, 6, 9, 10], a *Finsler space*  $F_n$ , determined by the symmetric and regular metric tensor  $g_{ij}(x^1, x^2, \dots, x^n, y^1, y^2, \dots, y^n)$ , determined by a function  $F(x, y) \equiv \frac{1}{2} L^2(x, y)$ :

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$$g_{ij}(x, y) = \frac{\partial^2 F(x, y)}{\partial y^i \partial y^j}.$$

The *Berwald connection* of  $F_n$  is introduced by the formulas:

$$G^h(x, y) = \frac{1}{2} g^{ij} \left( \frac{\partial^2 F(x, y)}{\partial y^j \partial x^k} y^k - \frac{\partial F(x, y)}{\partial y^j} \right),$$

$$G_i^h(x, y) = \frac{\partial G^h(x, y)}{\partial y^i}, \quad G_{ij}^h(x, y) = \frac{\partial^2 G^h(x, y)}{\partial y^i \partial y^j},$$

where  $g^{ij}$  are components of the matrix inverted to  $\|g_{ij}(x, y)\|$ .

For metric tensor of  $F_n$  we have the following formula:

$$g_{ij;k} = -2P_{ijk},$$

where  $P_{ijk} = C_{ijk;\alpha} y^\alpha$  is the Landsberg tensor, and  $C_{ijk} = \frac{1}{2} \frac{\partial g_{ij}}{\partial y^k}$  is the Cartan tensor. Hereafter the semicolon “;” denotes the *covariant derivative* with respect to the Berwald connection of the space  $F_n$ .

We consider two Finsler spaces  $F_n = (M_n, L)$  and  $\bar{F}_n = (M_n, \bar{L})$  on the same underlying manifold  $M_n$ . If any geodesic of a Finsler space  $F_n$  coincides with a geodesic of  $\bar{F}_n$  as a set of points and vice versa, the change  $L \rightarrow \bar{L}$  of the metrics is called *projective*, and the mapping  $F_n \rightarrow \bar{F}_n$  is a *geodesic mapping* between  $F_n$  and  $\bar{F}_n$ .

It is known that if a Finsler space  $F_n$  admits a geodesic mapping onto a Finsler space  $\bar{F}_n$  then in general with respect to the mapping system of coordinates  $x^1, x^2, \dots, x^n, y^1, y^2, \dots, y^n$  the objects of the Berwald connection of these spaces  $G^h(x, y)$  and  $\bar{G}^h(x, y)$  have the following relation [1, 2, 3, 5, 6, 9, 10]:

$$\bar{G}^h(x, y) = G^h(x, y) + \psi(x, y) y^h, \quad (1)$$

where  $\psi(x, y)$  is a positively homogeneous in  $y$  of degree one function.

The formula (1) equivalent to next formulas

$$\bar{G}_i^h(x, y) = G_i^h(x, y) + \delta_i^h \psi(x, y) + \psi_i(x, y) y^h,$$

$$\bar{G}_{ij}^h(x, y) = G_{ij}^h(x, y) + \delta_i^h \psi_j + \delta_j^h \psi_i + \psi_{ij}(x, y) y^h,$$

where

$$\psi_i = \frac{\partial \psi(x, y)}{\partial y^i}, \quad \psi_{ij} = \frac{\partial \psi_i(x, y)}{\partial y^j},$$

$\delta_i^h$  is the Kronecker symbol.

If  $\psi_i \neq 0$ , then a geodesic mapping is called *nontrivial*; otherwise it is said to be *trivial*.

From the above mentioned equations we obtain

$$\bar{G}_{ijk}^h(x, y) = G_{ijk}^h(x, y) + \delta_i^h \psi_{jk} + \delta_j^h \psi_{ik} + \delta_k^h \psi_{ij} + \psi_{ijk}(x, y) y^h,$$

where

$$\psi_{ijk} = \frac{\partial \psi_{ij}(x, y)}{\partial y^k}.$$

Finsler space is a *Berwald type* if  $G_{ijk}^h = 0$ , and Finsler space is a *weakly Berwald type* if  $G_{ij\alpha}^\alpha = G_{ij} = 0$ .

Hence  $\bar{G}_{ij} = G_{ij} + (n+1)\psi_{ij}$ , so if we consider the geodesic mapping between Berwald and weakly Berwald spaces then  $\psi_{ij} = 0$ .

It is well known, that every positive definit Berwald metric has common geodesic with some Riemannian metric [12]. The aim of this paper to show a new method, which in suitable for the investigation of geodesic mapping of Berwald spaces onto Riemannian spaces in any cases.

## 2. Geodesic Mapping of Weakly Berwald Spaces onto Riemannian Spaces

We know, that the Douglas tensor is invariant under geodesic mapping, that is  $D = \bar{D}$ , where

$$D_{ijk}^h = G_{ijk}^h - \frac{1}{(n+1)}(G_{ijk}y^h + G_{ij}\delta_k^h + G_{ik}\delta_j^h + G_{jk}\delta_i^h), \text{ and } G_{ijk} = \frac{\partial G_{ij}}{\partial y^k}.$$

The Douglas tensor vanishes in Riemannian spaces, so in the weakly Berwald space we get  $G_{ijk}^h = 0$ , where the weakly Berwald space has common geodesics with a Riemannian space. So we obtain the following

**Theorem 1.** *If a weakly Berwald space has common geodesics with a Riemannian space, then the weakly Berwald space is a Berwald space.*

Berwald constructed a tensor, which is analogically of Weyl tensor of projective curvature. This tensor is also invariant under a geodesic mappings of Finsler spaces.

The fundamental equation (1) of geodesic mapping between Finsler spaces  $F_n$  and  $\bar{F}_n$  equivalent to following formula:

$$\bar{g}_{ij;k} = 2\psi_k \bar{g}_{ij} + \psi_i \bar{g}_{jk} + \psi_j \bar{g}_{ik} + \bar{g}_{i\alpha} y^\alpha \psi_{jk} + \bar{g}_{j\alpha} y^\alpha \psi_{ik} - 2\bar{P}_{ijk}.$$

If case that space  $F_n$  and  $\bar{F}_n$  are weakly Berwald spaces this formula have

following simply form:

$$\bar{g}_{ij;k} = 2\psi_k \bar{g}_{ij} + \psi_i \bar{g}_{jk} + \psi_j \bar{g}_{ik} - 2\bar{P}_{ijk}. \quad (2)$$

In this case the functions  $\psi_i$  are independent of  $y$ , i.e.  $\psi_i = \psi_i(x)$ .

### 3. Geodesic Mappings of Berwald Spaces onto Riemannian Spaces

**Theorem 2.** *The Berwald space  $F_n$  admits a non-trivial geodesic mapping onto a Riemannian space  $\bar{V}_n$  with the metric tensor  $\bar{g}_{ij}(x)$  if and only if the following system of differential equations with covariant derivatives of Cauchy type has a solution with respect to the symmetric tensor  $\bar{g}_{ij}(x)$  ( $\det \|\bar{g}_{ij}(x)\| \neq 0$ ), the non-zero vector  $\psi_i(x)$  and the invariant  $\mu(x)$ :*

$$\begin{aligned} (a) \quad \bar{g}_{ij;k} &= 2\psi_k \bar{g}_{ij} + \psi_i \bar{g}_{jk} + \psi_j \bar{g}_{ik}; \\ (b) \quad n\psi_{i;j} &= n\psi_i \psi_j + \mu \bar{g}_{ij} - H_{ij} - \bar{g}_{i\alpha} \bar{g}^{\beta\gamma} H_{\beta\gamma}^\alpha - \frac{2}{n+1} H_{\alpha ij}^\alpha; \\ (c) \quad (n-1)\mu_{;i} &= 2(n-1)\psi_\alpha \bar{g}^{\beta\gamma} H_{\beta\gamma}^\alpha \\ &\quad + \psi_\alpha \bar{g}^{\alpha\beta} (5H_{\beta i} + \frac{6}{n+1} H_{\gamma\beta i}^\gamma - H_{i\beta}) \\ &\quad + \bar{g}^{\alpha\beta} (H_{\alpha\beta i;\gamma}^\gamma - H_{\alpha i;\beta} - \frac{2}{n+1} H_{\gamma\alpha i;\beta}^\gamma), \end{aligned} \quad (3)$$

where the semicolon denotes covariant derivative with respect to the space connection  $F_n$ ,  $\bar{g}^{ij}(x)$  are components of the matrix inverted to  $\|\bar{g}_{ij}(x)\|$ ,  $H_{ij}^h$  and  $H_{ij}$  are respectively curvature and Ricci tensors of the space  $F_n$ .

*Proof.* Suppose that  $F_n$  admits non trivial geodesic mapping onto  $\bar{V}_n$  with metric tensor  $\bar{g}_{ij}(x)$ . Then the connections  $F_n$  and  $\bar{V}_n$  have the relation (2) with  $\bar{P}_{ijk}^h = 0$ . If we consider the fundamental tensor  $\bar{g}_{ij}(x)$  of  $\bar{V}_n$ , then we get the conditions which are sufficient for  $F_n$  to admit non trivial geodesic mapping onto  $\bar{V}_n$ .

Let us consider integrability conditions of the equations (3a)

$$\bar{g}_{h\alpha} H_{ijk}^\alpha + \bar{g}_{i\alpha} H_{hjk}^\alpha = 2\bar{g}_{hi} \psi_{[jk]} + \bar{g}_j (h\psi_{i)k} - \bar{g}_k (h\psi_{i)j}, \quad (4)$$

where  $\psi_{ij} = \psi_{i;j} - \psi_i \psi_j$ ,  $[i j]$  and  $(i j)$  denotes alternation and symmetrization with respect to  $i$  and  $j$ , respective.

Transvecting (4) by  $\bar{g}^{ij}$ , we get  $\psi_{[jk]} = \frac{1}{n+1} H_{\alpha jk}^\alpha$ . Excluding  $\psi_{[jk]}$  from (4),

we obtain

$$\bar{g}_{\alpha(h}H_{i)jk}^{\alpha} - \frac{2}{n+1}\bar{g}_{hi}H_{\alpha jk}^{\alpha} = \bar{g}_{j(h}\psi_{i)k} - \bar{g}_{k(h}\psi_{i)j}. \tag{5}$$

After the transvecting (5) by  $g^{ik}$  we easily obtain the conditions (3b) with  $\mu \equiv \psi_{\alpha\beta}\bar{g}^{\alpha\beta}$ .

If we consider the equation  $\bar{g}^{ik}\bar{g}_{kj} = \delta_j^i$ , it is not difficult to show that the equations (3a) are equivalent to the relations

$$\bar{g}^{ij}{}_{;k} = -2\psi_k\bar{g}^{ij} - \delta_k^i\psi^j - \delta_k^j\psi^i, \tag{6}$$

where  $\psi^i \equiv \psi_{\alpha}\bar{g}^{\alpha i}$ .

We covariantly differentiate the conditions (3b) and then alternate the result with respect to the indices  $j$  and  $k$  taking into account (3a), (3b), (6) and transvecting by  $\bar{g}^{ik}$ , and finally we get equations (3c).

The theorem has been proved. □

From Theorem 2 we may conclude that the set of all Riemannian spaces  $\bar{V}_n$ , the given Berwald space  $F_n$  admits non trivial geodesic mapping onto  $\bar{V}_n$ , is dependent on  $r \leq r_0 = (n+1)(n+2)/2$  parameters.

Finding of all the solutions of (3) requires a consideration of their integrability conditions and differential extensions, which form a set of algebraic equations with respect to the unknown functions  $\bar{g}_{ij}$ ,  $\psi_i$  and  $\mu$  with coefficients from  $F_n$ . But this set is not linear and its solution is certainly difficult.

The results under discussion are generalisations of analogous theorems of N.S. Sinyukov [11] and J. Mikeš, V. Berezovski [7] for the geodesic mappings of Riemannian spaces.

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