

# ALMOST GEODESIC MAPPINGS OF TYPE $\pi_1^*$ OF SPACES WITH AFFINE CONNECTION

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**Summary.** We consider almost geodesic mappings  $\pi_1^*$  of spaces with affine connections. This mappings are a special case of first type almost geodesic mappings. We have found the objects which are invariants of the mappings  $\pi_1^*$ . The fundamental equations of these mappings are in Cauchy form. We study  $\pi_1^*$  mappings of constant curvature spaces.

## 1 INTRODUCTION

In the theory of geodesic mappings and their generalizations many basic results were formulated as a system of differential equations in Cauchy form, see [1–14]. For almost geodesic mappings  $\pi_1$  a similar result for special Ricci-Codazzi Riemannian spaces is formulated in Sinyukov monograph [1]. This result was generalized for Ricci-Codazzi spaces with affine connection and for Riemannian spaces in [15]. For  $\pi_1^*$  mappings of general symmetric spaces with affine connection the system of differential equations in the Cauchy form were found in works [16].

This paper is devoted to detailed study of  $\pi_1^*$  mappings which are characterized by the general equations in the Cauchy form. This result is significant because the equations in this form have established methods of solution.

The concept of almost geodesic mappings of type  $\pi_1^*$  of spaces with affine torsion-free connections was first introduced in [16]. These mappings are a special case of type  $\pi_1$  almost geodesic mappings which were introduced by N. S. Sinyukov in [1].

The paper is devoted to study the general properties of  $\pi_1^*$  mappings. In particular, we have obtained the objects which are invariant under the mappings. Also  $\pi_1^*$  mappings of spaces of constant curvature and affine spaces were studied.

Let us recall the basic conceptions of the almost geodesic mappings theory presented in [1].

A curve defined in a space with an affine connection  $A_n$  is called *almost geodesic* if there exists a two-dimensional plane element parallel along the curve (relative to the affine connection) such that for any tangent vector of the curve its parallel translation along the curve belongs to the plane element.

A diffeomorphism  $f$  between spaces with affine connection  $A$  and  $\bar{A}_n$  is called *almost geodesic mapping* if any geodesic curve of  $A$  is mapped under  $f$  onto an almost geodesic curve in  $\bar{A}$ .

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In order that a mapping of a space  $A_n$  onto a space  $\bar{A}_n$  be almost geodesic it is necessary and sufficient that in a common coordinate system  $x \equiv (x^1, x^2, \dots, x^n)$  which both spaces are referred to, the deformation tensor of the mapping  $P_{ij}^h(x) \equiv \bar{\Gamma}_{ij}^h(x) - \Gamma_{ij}^h(x)$  must satisfy the conditions

$$A_{\alpha\beta\gamma}^h \lambda^\alpha \lambda^\beta \lambda^\gamma = a \cdot P_{\alpha\beta}^h \lambda^\alpha \lambda^\beta + b \cdot \lambda^h,$$

where  $A_{ijk}^h \equiv P_{ij,k}^h + P_{ij}^\alpha P_{\alpha k}^h$ ,  $\Gamma_{ij}^h(x)$  ( $\bar{\Gamma}_{ij}^h(x)$ ) are the components of the affine connection of the space  $A_n$  ( $\bar{A}_n$ ),  $''$  denotes covariant derivative with respect to the connection of the space  $A_n$ ,  $\lambda^h$  is an arbitrary vector,  $a$  and  $b$  are certain functions of variables  $x^h$  and  $\lambda^h$ .

Three types of almost geodesic mappings were specified, namely  $\pi_1$ ,  $\pi_2$ ,  $\pi_3$ . We have proved that for  $n > 5$  other types of almost geodesic mappings except  $\pi_1$ ,  $\pi_2$ , and  $\pi_3$  do not exist [17].

Almost geodesic mappings of  $\pi_1$  type are characterized by the following conditions for the deformation tensor:

$$A_{(ijk)}^h = \delta_{(i}^h a_{jk)} + b_{(i} P_{jk)}^h,$$

where  $a_{ij}$  is a certain symmetric tensor,  $b_i$  a certain covector,  $\delta_i^h$  are the Kronecker delta,  $(ijk)$  denotes an operation called symmetrization without division with respect to the indices  $i, j$  and  $k$ .

Unlike mappings of the type  $\pi_1$ , the study of mappings  $\pi_2$  and  $\pi_3$  are devoted by a lot of papers (See e.g. [1,2]). It stems from the fact that the main equations of these mappings are much more sophisticated than equations of other ones. Hence the paper is devoted to a special case of mappings  $\pi_1$ , which does not degenerate into  $\pi_2$ ,  $\pi_3$  or geodesic mappings.

## 2 ALMOST GEODESIC MAPPINGS OF THE $\pi_1^*$ TYPE

Let a mapping of  $A_n$  onto  $\bar{A}_n$  satisfy the conditions [16]:

$$P_{ij,k}^h + P_{ij}^\alpha P_{\alpha k}^h = a_{ij} \delta_k^h, \quad (1)$$

where  $a_{ij}$  is a certain symmetric tensor.

These mappings are a special case of almost geodesic mappings of the  $\pi_1$  type. From now on, that mappings will be denoted by  $\pi_1^*$ . Let us consider (1) as a system of differential equations of Cauchy type with respect to the deformation tensor  $P_{ij}^h$  and find their integrability conditions. To this end, differentiate covariantly (1) with respect to  $x^m$  in  $A_n$ , then alternate it in  $k$  and  $m$ .

Contracting the integrability conditions of the equations (1) for  $h$  and  $m$ , we get

$$(n-1)a_{ij,k} = P_{ij}^\alpha R_{\alpha k} - P_{\alpha(i}^\beta R_{j)\beta k}^\alpha - (n-1)P_{ij}^\alpha a_{\alpha k}, \quad (2)$$

where  $R_{ijk}^h$  is the Riemann tensor of the space  $A_n$ ,  $R_{ij}^\alpha \equiv R_{ij\alpha}^\alpha$  is the Ricci tensor.

Obviously, in the space  $A_n$  the equations (1) and (2) form a closed system of PDEs of Cauchy type with respect to the functions  $P_{ij}^h(x)$  and  $a_{ij}(x)$ . The functions must also satisfy the algebraic conditions

$$P_{ij}^h(x) = P_{ji}^h(x), \quad a_{ij}(x) = a_{ji}(x). \quad (3)$$

Hence we have proved the theorem.

**Theorem 1** *In order that a space  $A_n$  with an affine connection admits a canonical almost geodesic mapping of type  $\pi_1^*$  onto another space  $\bar{A}_n$  with an affine connection, it is necessary and sufficient that the mixed system of differential equations of Cauchy type in covariant derivatives (1), (2), (3) has a solution with respect to the unknown functions  $P_{ij}^h(x)$  and  $a_{ij}(x)$ .*

Let us note that Theorem 1 holds for  $A_n \in C^1$  ( $\Gamma_{ij}^h(x) \in C^1$ ), i.e. objects of affine connection  $\Gamma$  are differentiable. In this case, if  $A_n \in C^r$  ( $r \geq 1$ ) then  $\bar{A}_n \in C^r$ . It follows from the fact, that the solution  $P_{ij}^h(x) \in C^r$  and  $a_{ij}(x) \in C^{r-1}$ .

The integrability conditions of the system are

$$-P_{ij}^\alpha R_{\alpha km}^h + P_{\alpha(i}^h R_{j)km}^\alpha = \frac{1}{n-1} [(P_{ij}^\alpha R_{\alpha m}^h - P_{\alpha(i}^\beta R_{j)m\beta}^\alpha) \delta_k^h - (P_{ij}^\alpha R_{\alpha k}^h - P_{\alpha(i}^\beta R_{j)k\beta}^\alpha) \delta_m^h],$$

where  $[ij]$  denotes the alternation with respect to the mentioned indices.

### 3 INVARIANT OBJECTS UNDER $\pi_1^*$ MAPPINGS

It is known [1], that if  $P_{ij}^h$  is a deformation tensor, then the Riemann tensors  $R_{ijk}^h$  and  $\bar{R}_{ijk}^h$  of the spaces  $A_n$  and  $\bar{A}_n$  are related to each other by the equations

$$\bar{R}_{ijk}^h = R_{ijk}^h + P_{i[k,j]}^h + P_{i[k}^\alpha P_{j]\alpha}^h. \quad (4)$$

Using the formulas (1) and (4), we get

$$* \bar{W}_{ijk}^h = * W_{ijk}^h, \quad (5)$$

where

$$* W_{ijk}^h \equiv R_{ijk}^h - \frac{1}{n-1} R_{i[j}^h \delta_{k]}^h, \quad * \bar{W}_{ijk}^h \equiv \bar{R}_{ijk}^h - \frac{1}{n-1} \bar{R}_{i[j}^h \delta_{k]}^h. \quad (6)$$

Obviously,  $* W_{ijk}^h$  is a tensor of type (1,3) in the space  $A_n$ , and  $* \bar{W}_{ijk}^h$  is a tensor of the same type in the space  $\bar{A}_n$ . From the relations (5) it follows that the tensor is invariant under almost geodesic mappings  $\pi_1^*$ .

Contracting (5) for  $h$  and  $i$ , it is easy to see that it holds

$$W_{ij} = \bar{W}_{ij}, \quad (7)$$

where

$$W_{ij} \equiv R_{[ij]}, \quad \bar{W}_{ij} \equiv \bar{R}_{[ij]}. \quad (8)$$

Taking account of (7), the formulas (5) are expressible in the form

$$\bar{W}_{ijk}^h = W_{ijk}^h, \quad (9)$$

where  $W_{ijk}^h$  and  $\bar{W}_{ijk}^h$  are the Weyl tensors of projective curvature of the spaces  $A_n$  and  $\bar{A}_n$  respectively.

Finally we obtained the theorem.

**Theorem 2** *The Weyl tensor of projective curvature  $W_{ijk}^h$ , and also the tensors  $*W_{ijk}^h$  and  $W_{ij}$  defined by the formulas (6) and (8) as geometric objects of spaces with affine connections are invariant under almost geodesic mappings of type  $\pi_1^*$ .*

#### 4 MAPPINGS $\pi_1^*$ OF EQUIAFFINE AND PROJECTIVE EUCLIDEAN SPACES

From the Theorem 2 we obtain the next one.

**Theorem 3** *If a projective Euclidean space admits an almost geodesic mappings of type  $\pi_1^*$  onto  $\bar{A}_n$ , then  $\bar{A}_n$  itself is also a projective Euclidean space.*

**Theorem 4** *If an equiaffine space admits an almost geodesic mappings of type  $\pi_1^*$  onto  $\bar{A}_n$ , then  $\bar{A}_n$  itself is also an equiaffine space.*

*Proof.* Obviously, the proof of Theorem 3 and 4 follows from the facts that the Weyl tensor of projective curvature vanishes in a projective Euclidean space, and for an equiaffine space the condition  $W_{ij} = 0$  holds identically, respectively.

Hence because of Theorem 2 the above mentioned tensors vanish in the space  $\bar{A}_n$ . This means that  $\bar{A}_n$  is a projective Euclidean and equiaffine space, respectively.

Thus from Theorem 3 and 4 projective Euclidean and equiaffine spaces form closed classes with respect to mappings of type  $\pi_1^*$ .

It is easy to see that the Riemann tensor is preserved under mappings  $\pi_1^*$  if and only if the tensor  $a_{ij}$  vanishes identically. In this case the main equations of the mappings become

$$P_{ij,k}^h = -P_{ij}^\alpha P_{\alpha k}^h. \quad (10)$$

In an affine space the equations (10) are completely integrable. Consequently, a solution of the equations is determined by arbitrary initial values of  $P_{ij}^h(x_0)$ . If the initial values satisfy the condition  $P_{ij}^h(x_0) \neq \delta_{(i}^h \psi_{j)}(x_0)$ , then the constructed solution determines the mapping of an affine space  $A_n$  onto another affine space  $\bar{A}_n$ , and the mapping is different from a geodesic one.

Hence we obtain the theorem.

**Theorem 5** *There is a mapping  $\pi_1^*$  of affine space onto itself such that all straight lines are mapped onto plain curves, and not all the curves are straight lines.*

Moreover, since in affine spaces the integrability conditions (2) of the equations (1) are satisfied identically, the equations (1) are completely integrable.

Let us prove the theorem.

**Theorem 6** *Riemannian spaces  $V_n$  of non-zero constant curvature admit non-geodesic mappings  $\pi_1^*$ , which are also almost geodesic mappings of type  $\pi_3$ . The quadratic complex of geodesics is preserved under the mappings.*

*Proof.* Let  $V_n$  be a Riemannian spaces with non-zero constant curvature  $R$  which admits non-geodesic mappings  $\pi_1^*$ . The integrability conditions are expressible in the form

$$K(P_{k(i)g_{jl}}^h - P_{l(i)g_{jk}}^h) + \delta_l^h B_{ijk} - \delta_k^h B_{ijl} = 0, \quad (11)$$

where  $B_{ijk} \equiv a_{ij,k} + P_{ij}^\alpha(a_{\alpha k} + K g_{\alpha k})$ ,  $g_{ij}$  is the metric tensor of the space  $V_n$ .

Let  $\epsilon^h$  be a vector such that  $\epsilon^\alpha \epsilon^\beta g_{\alpha\beta} = \pm 1$ . Transvecting (11) with  $\epsilon^j \epsilon^l$  and then symmetrizing it in  $i$  and  $k$ , we find

$$P_{ik}^h = \xi^h g_{ik} + \epsilon^h b_{ik} + \delta_{(i}^h \psi_{j)}, \quad (12)$$

where  $\xi^h$ ,  $\psi_j$  are some vectors,  $b_{ik}$  is some symmetric tensor. From (12) it follows that the relation (11) becomes

$$\epsilon^h (b_{k(i)g_{jl}} - b_{l(i)g_{jk}}) + \delta_{[l}^h b_{k]ij} + \delta_{(i}^h g_{j)[l} \psi_{k]} = 0, \quad (13)$$

where  $b_{ijk}$  is some tensor.

Transvecting (13) with  $\epsilon^l$ , we get

$$\delta_i^h (g_{j\alpha} \epsilon^\alpha \psi_k - g_{jk} \epsilon^\alpha \psi_\alpha) + \epsilon^h b_{1ijk} + \delta_j^h b_{2ik} + \delta_k^h b_{3ij} = 0, \quad (14)$$

where  $b_{1ijk}$ ,  $b_{2ik}$ ,  $b_{3ij}$  are some tensors.

Suppose that  $\psi_i \neq 0$ . Then  $g_{j\alpha} \epsilon^\alpha \psi_k - g_{jk} \epsilon^\alpha \psi_\alpha \neq 0$  and consequently there exist vectors  $a^j$  and  $b^k$  such that  $a^j b^k (g_{j\alpha} \epsilon^\alpha \psi_k - g_{jk} \epsilon^\alpha \psi_\alpha) \neq 0$ . Transvecting (13) with  $a^j b^k$ , we obtain a relation which is contrary to the assumption that  $n > 3$ . Hence  $\psi_i = 0$ . The formulas (14) can be simplified and we can show by a similar method that  $b_{kij} = 0$ . Then (13) becomes  $b_{k(i)g_{jl}} - b_{l(i)g_{jk}} = 0$ . Transvecting the latter with  $g^{jl}$ , we find that  $b_{ki} = \frac{b_{\alpha\beta} g^{\alpha\beta}}{n} g_{ki}$ . We have from (12) by direct calculation

$$P_{ij}^h = P^h g_{ij}, \quad (15)$$

where  $P^h$  is some vector. Hence the mapping is  $f$ -planar. Consequently, according to [1,2], such mapping is almost geodesic mapping of type  $\pi_3$ . And in [17] the authors have proved that the mappings  $\pi_1 \cap \pi_3$  preserves the quadratic complex of geodesics [18].

Substituting (15) in (1), we have

$$P_{,k}^h + P^h P_k = \alpha \delta_k^h,$$

where  $\alpha$  is some invariant,  $P_k = P^l g_{lk}$ .

Vector fields satisfying these conditions are referred to as *concircular vector fields*. One knows that concircular vector fields always exist in spaces of constant curvature.

## 5 EXAMPLES OF ALMOST GEODESIC MAPPINGS $\pi_1^*$

We shall give an example of an almost geodesic mapping  $\pi_1^*$  of a flat space  $A_n$  onto another flat space  $\bar{A}_n$ .

Let  $x^1, x^2, \dots, x^n$  and  $\bar{x}^1, \bar{x}^2, \dots, \bar{x}^n$  be affine coordinate systems in the spaces  $A_n$  and  $\bar{A}_n$  respectively. A point mapping

$$\bar{x}^h = \frac{1}{2} C_\alpha^h (x^\alpha - C^\alpha)^2 + x_0^h, \quad (16)$$

where  $C_i^h, C^h, x_0^h$  are some constant,  $\det \| C_i^h \| \neq 0$ , defines the almost geodesic mapping  $\pi_1^*$  of the space  $A_n$  onto the space  $\bar{A}_n$ .

By direct calculation it is readily shown that the components of the deformation tensor  $P_{ij}^h$  in the coordinate system  $x^1, x^2, \dots, x^n$  are given

$$P_{ii}^i = \frac{1}{x^i - C^i} \quad i = \overline{1, n},$$

all the other components being zero.

Obviously, the tensor satisfies the equation (10). Note that the mapping is different from mappings of types  $\pi_2$  and  $\pi_3$ .

Straight lines which are defined in the space  $A_n$  by the equation  $x^h = a^h + b^h t$  ( $t$  is a parameter along a line) are mapped into parabolas in the space  $\bar{A}_n$ . The parabolas are defined by the equations

$$\bar{x}^h = F^h + D^h t + E^h t^2,$$

where  $F^h = \frac{1}{2} C_\alpha^h (a^\alpha - C^\alpha)^2$ ,  $D^h = C_\alpha^h (a^\alpha - C^\alpha) b^\alpha$ ,  $E^h = \frac{1}{2} C_\alpha^h (b^\alpha)^2$ .

The exceptions are the straight lines through the point  $M(C^1, C^2, \dots, C^n)$ . By (16) the lines are mapped into straight lines too.

Finally we note that the formulas (16) generate a family of almost geodesic mappings  $\pi_1$  of a flat space if the parameters  $C_i^h, C^h$  and  $x_0^h$  are understood as continuous values.

## 6 CONCLUSION

Out of the three types of almost geodesic mappings of spaces with affine connection, distinguished by N.S. Sinyukov, the least studied are almost geodesic mappings of the first type. The equations that characterize them are very complex. Therefore, the results obtained for mappings  $\pi_1^*$ , including for their particular cases, are very relevant and are of theoretical value from the geometrical point of view. At the same time, they can be used in the theory of relativity and theoretical mechanics.

Almost geodesic mappings are a natural generalization of geodesic mappings. The basic equations of geodesic mappings of spaces with affine connection cannot be reduced to closed systems of equations in covariant derivatives of Cauchy type, since the general solution depends on  $n$  arbitrary functions.

We have singled out a special case of almost geodesic mappings of the first type, denoted by  $\pi_1^*$ , the basic equations of which are reduced to a closed system of equations in covariant derivatives of the Cauchy type. This result is very important, since (since geodesic mappings are a special case of almost geodesic mappings) the basic equations of the first type of spaces with affine connection are not reducible to closed systems of equations in covariant derivatives of Cauchy type.

For the mappings  $\pi_1^*$  geometric objects of tensor nature are found that are invariant under such mappings. It turns out that the Weyl tensor is invariant not only with respect to geodesic mappings, but also with respect to more general mappings.

In the article it is proved that projective-Euclidean and equiaffine spaces form closed classes with respect to mappings  $\pi_1^*$ .

From geometrical point of view, an interesting is a special case of mappings  $\pi_1^*$ , which we have distinguished, where the Riemann tensor is invariant. In this case, the basic equations of such mappings in flat space are completely integrable. An example of mappings  $\pi_1^*$  of flat space onto flat space is given.

In wpresented paper, it is of interest to study the integrability conditions and their differential extensions of the obtained equations in covariant derivatives of the Cauchy type that characterize the mappings  $\pi_1^*$  of spaces with affine connection.

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